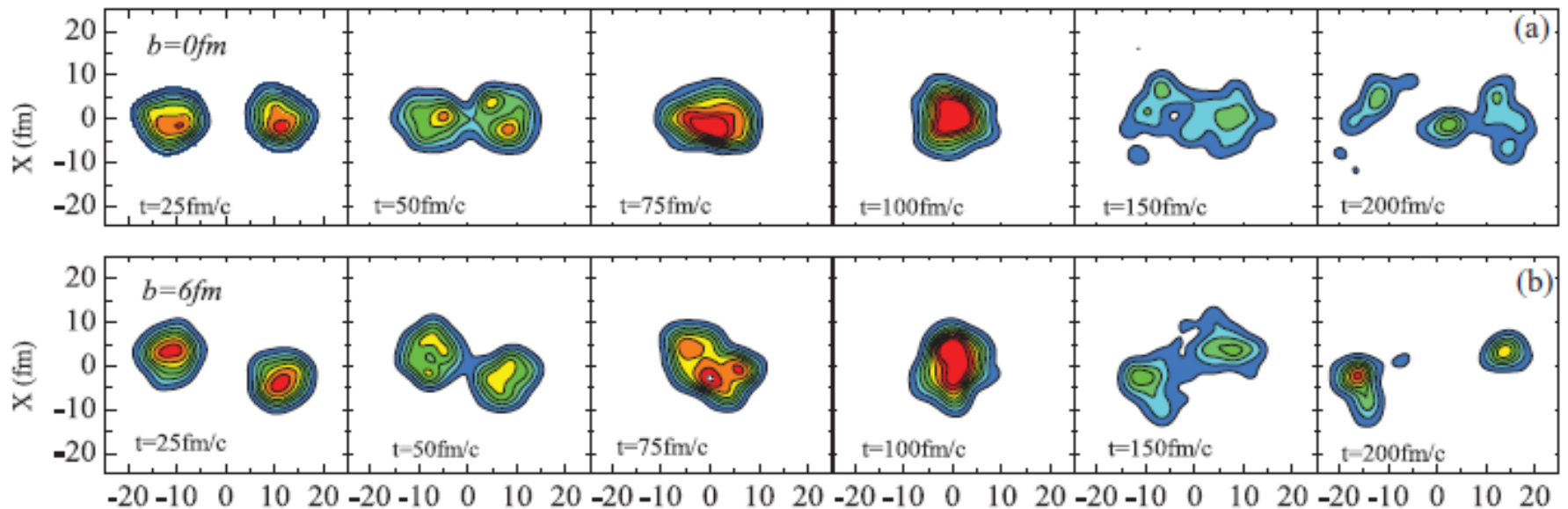


1st Symposium on Intermediate-energy Heavy Ion Collisions (iHIC2018)

Bayesian inference on Skyrme EDF from HICs

Yingxun Zhang (张英逊)



Outline

1, Nuclear equation of state, Skyrme interaction and its constraints

2, preliminary constraints on Skyrme parameter sets from isospin sensitive HIC observables

- Sensitivity study of model parameter x and observables O in transport model
- CI-DR(n/p) at 120 A MeV, R_{diff} at 50 A MeV, and $R_i(y)$ at 35, 50 A MeV

3, Summary and outlook

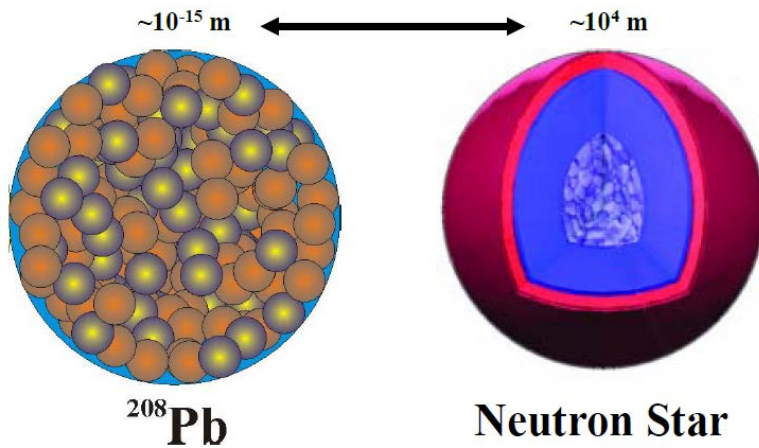
Isospin asymmetric Equation of State

$$E(\rho, \delta) = E(\rho, \delta = 0) + S(\rho)\delta^2 + O(\delta^4)$$

It is a fundamental properties of nuclear matter, and is very important for understanding

- *properties of nuclear structure*
- *properties of neutron star*
- *properties of heavy ion reaction mechanism*

Symmetry energy on vastly differing length scales



extrapolation from ^{208}Pb radius to n-star radius

C.J. Horowitz, J. Piekarewicz, Phys. Rev. Lett. 86 (2001) 5647

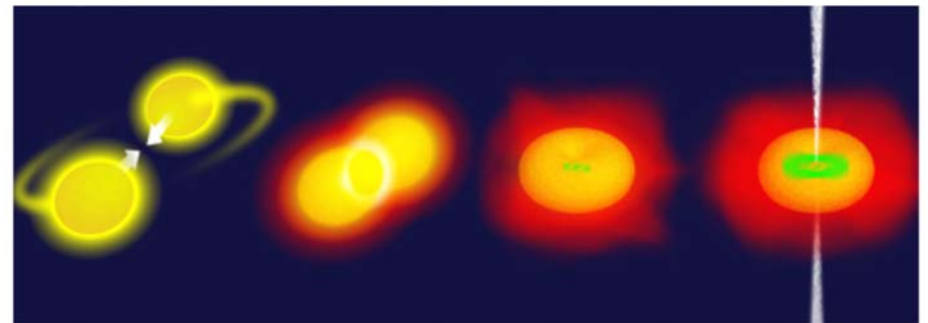
PRL 119, 161101 (2017)

Selected for a Viewpoint in *Physics*
PHYSICAL REVIEW LETTERS

week ending
20 OCTOBER 2017

GW170817: Observation of Gravitational Waves from a Binary Neutron Star Inspiral

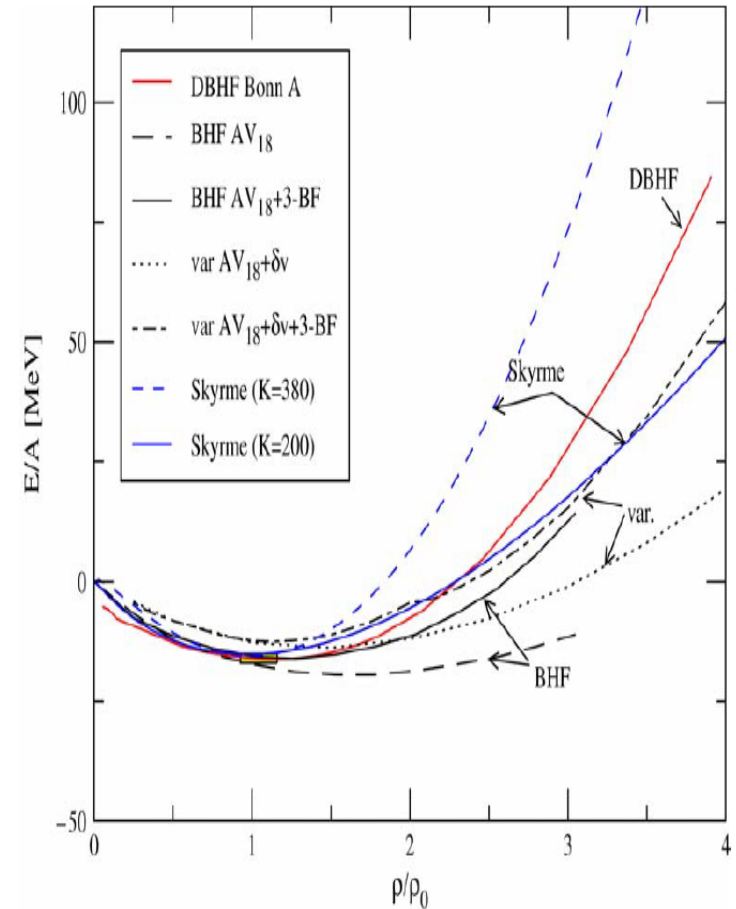
B. P. Abbott *et al.*^{*}
(LIGO Scientific Collaboration and Virgo Collaboration)
(Received 26 September 2017; revised manuscript received 2 October 2017; published 16 October 2017)
On August 17, 2017 at 12:41:04 UTC the Advanced LIGO and Advanced Virgo gravitational-wave



Theoretical predictions on the properties of nuclear matter

- Effective field theory approaches (Based on Chiral perturbation theory,
- Ab initio approaches (Based on the high precision free space nucleon-nucleon interaction) DBHF, SCGF, QMC
- Phenomenological density functional (Based on Gogny or **Skyrme** force or RMF,

C. Fuchs / Progress in Particle and Nuclear Physics 56 (2006) 1–103



• Skyrme interaction (1956)

E. Chabanat et al./Nuclear Physics A 627 (1997) 710–746

$$\begin{aligned}
 V(\mathbf{r}_1, \mathbf{r}_2) = & t_0 (1 + x_0 P_\sigma) \delta(\mathbf{r}) \\
 & + \frac{1}{2} t_1 (1 + x_1 P_\sigma) \left[\mathbf{P}'^2 \delta(\mathbf{r}) + \delta(\mathbf{r}) \mathbf{P}^2 \right] \\
 & + t_2 (1 + x_2 P_\sigma) \mathbf{P}' \cdot \delta(\mathbf{r}) \mathbf{P} \\
 & + \frac{1}{6} t_3 (1 + x_3 P_\sigma) [\rho(\mathbf{R})]^\sigma \delta(\mathbf{r}) \\
 & + iW_0 \boldsymbol{\sigma} \cdot [\mathbf{P}' \times \delta(\mathbf{r}) \mathbf{P}]
 \end{aligned}$$

central term

Effective Skyrme

non-local terms

density-dependent term

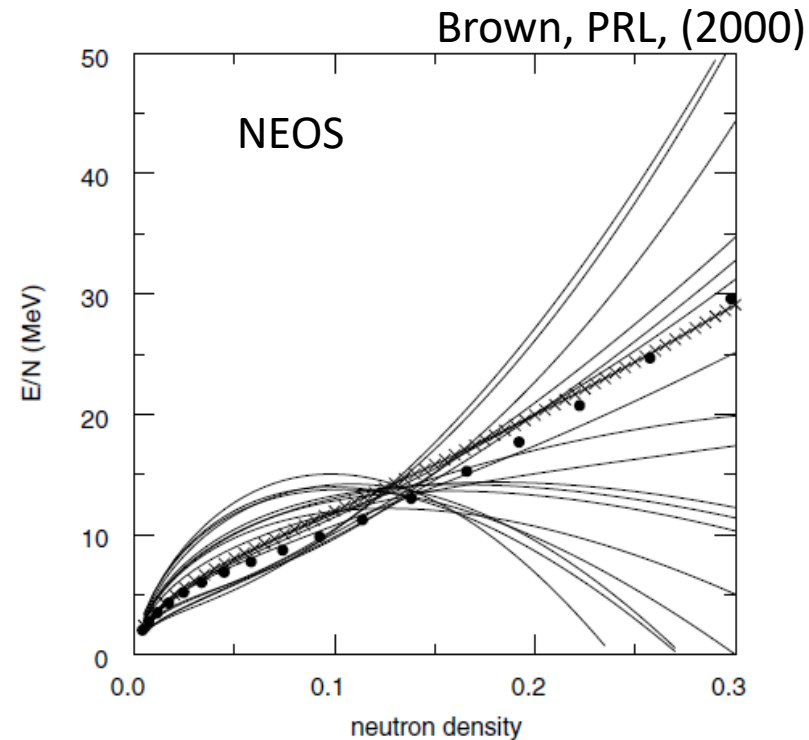
spin-orbit term .

• **simple**, effectively taken into account the complicated correlations

• parameter are determined by fitting, ρ_0 , K_0 , m_s^* , κ , $mass$,

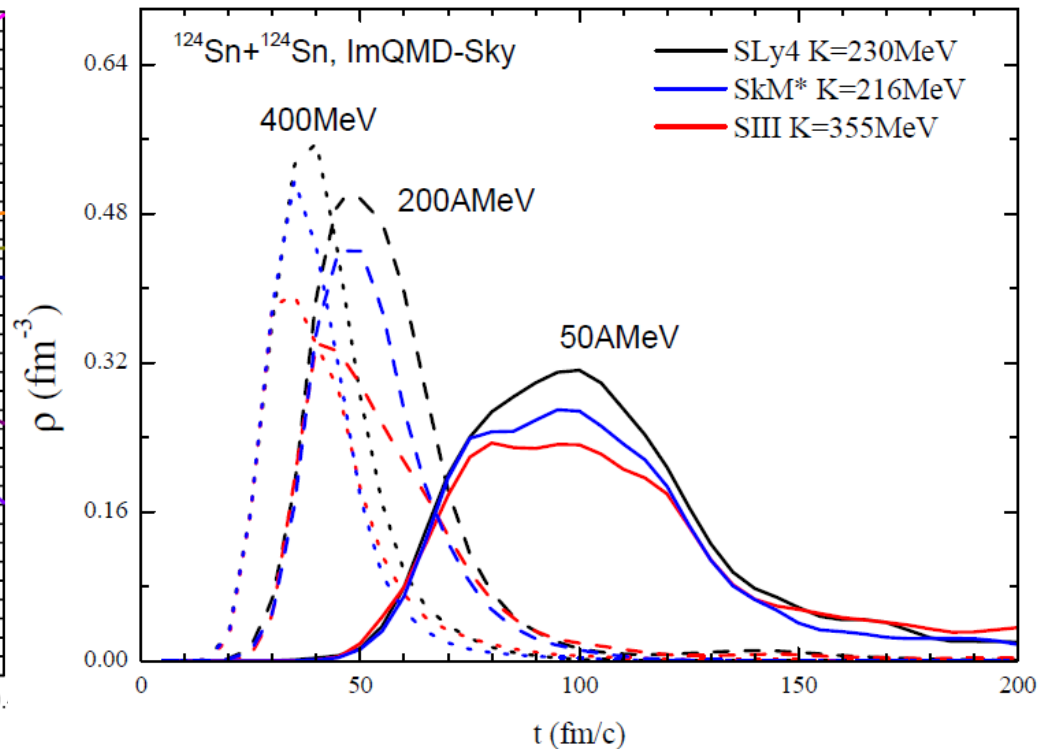
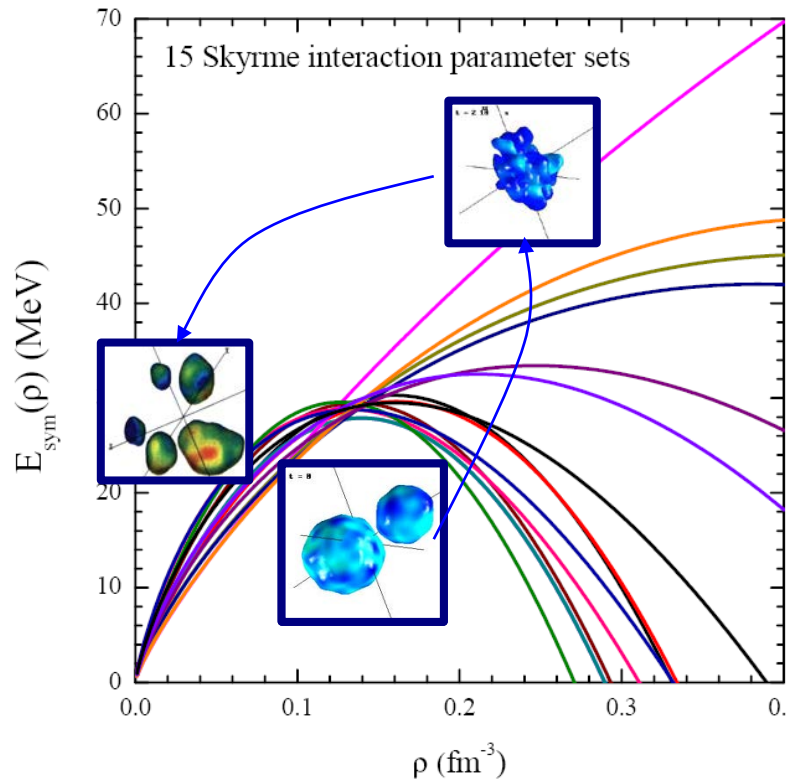
• **Many applications in nuclear structure studies**

• *parameterization is not unique and there exists >240 sets (drawback)*



Can we learn the information of equation of state away from normal density?

Yes, we can! Heavy ion collisions + Transport model calculations



Input of transport models:

nucleonic potential or interaction parameter (\rightarrow EoS)

A, BUU type: $f(\mathbf{r}, \mathbf{p}, t)$ one body phase space density

$$\frac{\partial f}{\partial t} + \mathbf{v} \cdot \nabla_{\mathbf{r}} f - \nabla_{\mathbf{r}} U \cdot \nabla_{\mathbf{p}} f$$

Mean field

$$f(\mathbf{r}, \mathbf{p}) \cong \frac{1}{\tilde{N}} \sum_{i=1}^{N\tilde{N}} \delta(\mathbf{r} - \mathbf{r}_i) \delta(\mathbf{p} - \mathbf{p}_i)$$

$$= -\frac{1}{(2\pi)^6} \int d^3 p_2 d^3 p_{2'} d\Omega \frac{d\sigma}{d\Omega} v_{12} \\ \times \{ [f f_2 (1 - f_{1'}) (1 - f_{2'})] - [f_{1'} f_{2'} (1 - f) (1 - f_2)] \\ \times (2\pi)^3 \delta^3(\mathbf{p} + \mathbf{p}_2 - \mathbf{p}_{1'} - \mathbf{p}_{2'}) \}$$

Two-body collision: occurs between test part.

Solved with test particle methods

B, QMD type: solve N-body equation of motion

$$\phi_i(\mathbf{r}_i, t) = \frac{1}{(2\pi\sigma_r^2)^{3/4}} \exp\left[-\frac{(\mathbf{r} - \mathbf{r}_i)^2}{4\sigma_r^2} + \frac{i\mathbf{p}_i \cdot \mathbf{r}}{\hbar}\right] \quad \text{nucleon}$$

$$\dot{\mathbf{r}}_i = \frac{\mathbf{p}_i}{m} + \nabla_{\mathbf{p}_i} \sum_j \langle V_{ij} \rangle = \nabla_{\mathbf{p}_i} \sum_j \langle H \rangle$$

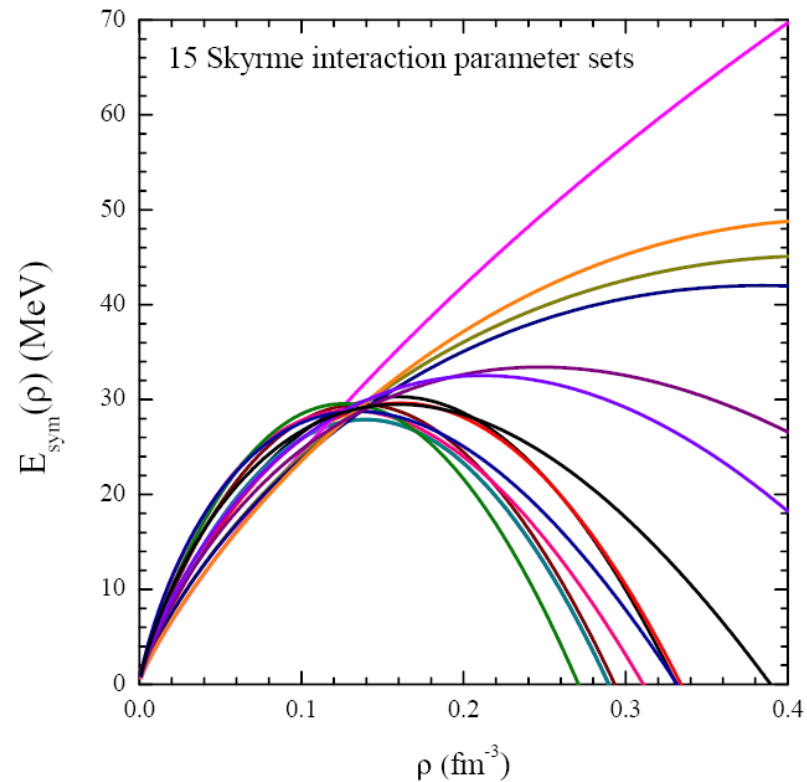
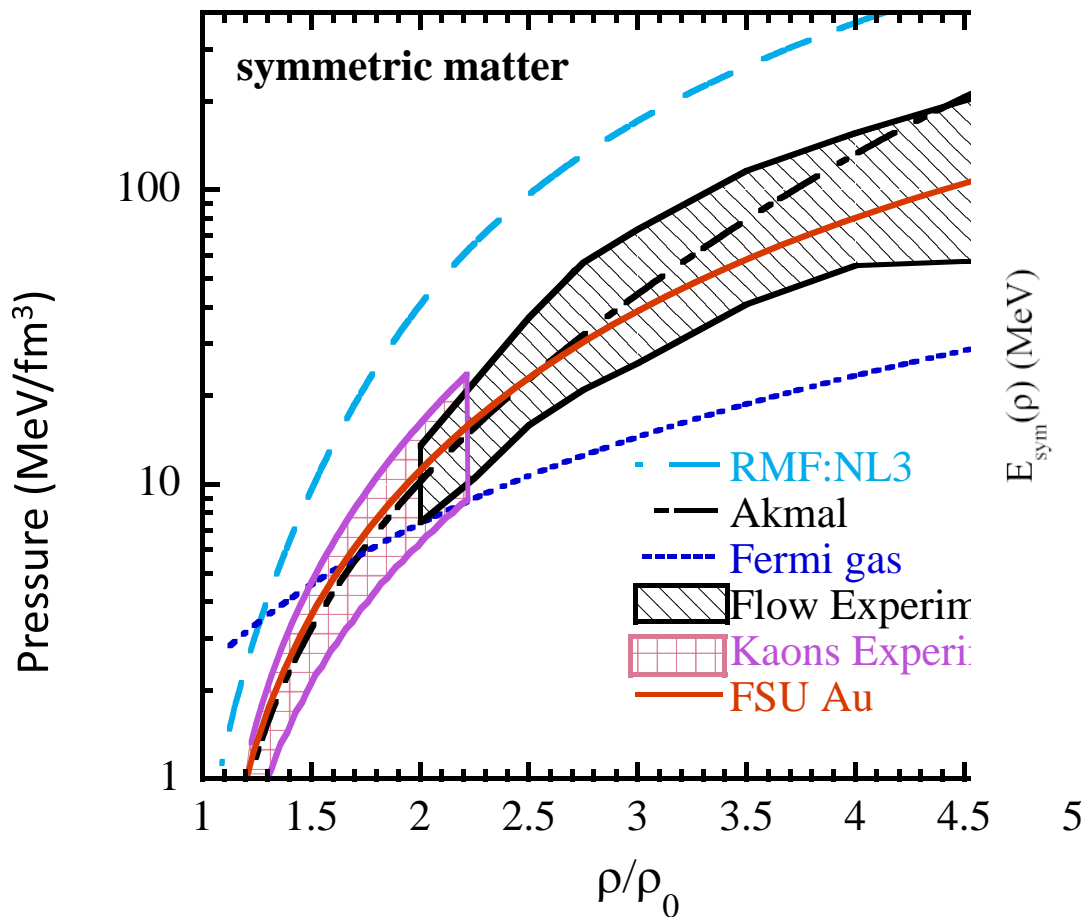
$$\dot{\mathbf{p}}_i = -\nabla_{\mathbf{r}_i} \sum_j \langle V_{ij} \rangle = -\nabla_{\mathbf{r}_i} \sum_j \langle H \rangle$$

Two body collision: occurs between nucleons
Rearrange whole nucleon \rightarrow large fluctuation

- symmetric part is relatively well constrained

$$E/A(\rho, \delta) = E/A(\rho, 0) + \delta^2 \cdot S(\rho) ; \delta = (\rho_n - \rho_p) / (\rho_n + \rho_p) = (N-Z)/A$$

Danielewicz, Lacey, Lynch, Science 298,1592 (2002)

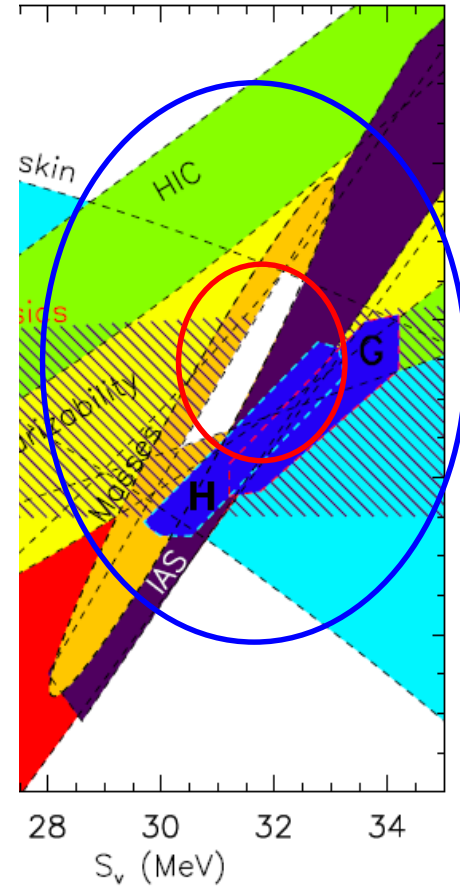
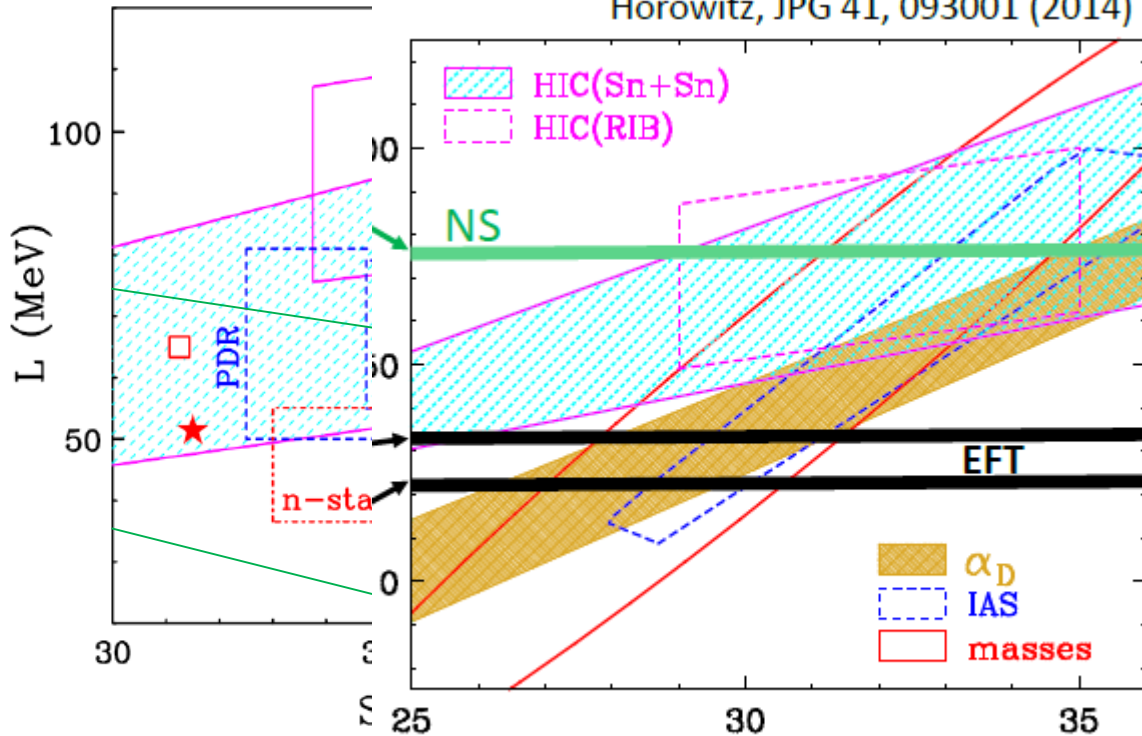


Progress on the constraints of symmetry energy

M.B.Tsang, Yingxun Zhang, et.al., PRL2009

Lattimer, EPJA 50 (2014) 40

Horowitz, JPG 41, 093001 (2014)



$$S(\rho) = S_0 + \frac{L}{3} \left(\frac{\rho - \rho_0}{\rho_0} \right) + \frac{S_0}{18} \left(\frac{\rho - \rho_0}{\rho_0} \right)^2 + \dots,$$

Consensus on symmetry energy have been obtained at subsaturation density.

Uncertainties on the constraints still need to be understand and constraints on symmetry energy need to be improved! ($L \pm \Delta L$)

Parameter correlations in $S(\rho)$: $\{L, m_s^*, f_I\}$ in Skyrme-HF

- Density dependent of symmetry energy from SHF

$$S(\rho) = \frac{1}{3} \frac{\hbar^2}{2m} \left(\frac{3\pi^2}{2}\right)^{2/3} \rho^{2/3} - \frac{1}{8} t_0 (2x_0 + 1) \rho - \frac{1}{48} t_3 (2x_3 + 1) \rho^{\sigma+1} - \frac{1}{24} \left(\frac{3\pi^2}{2}\right)^{2/3} (3\Theta_v - 2\Theta_s) \rho^{5/3} \quad (C8)$$

$$S(\rho) = \frac{1}{3} \epsilon_F \rho^{2/3} + A_{sym} \rho + B_{sym} \rho^{\sigma+1} + C_{sym} (m_s^*, m_v^*) \rho^{5/3} \quad (C11)$$

- Density dependent of symmetry depends not only on **density**, **effective mass splitting** but also **isoscalar effective mass**

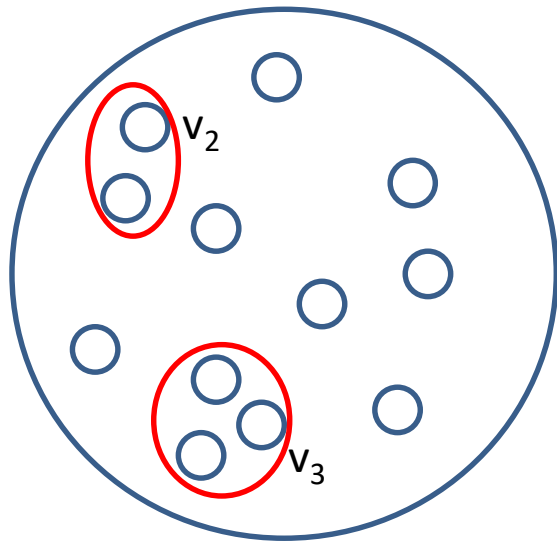
$$f_I = \frac{1}{2\delta} \left(\frac{m}{m_n^*} - \frac{m}{m_p^*} \right) = \frac{m}{k} \frac{\partial((U_n - U_p)/2\delta)}{\partial k} = \frac{\partial U_{sym}}{\partial E_k}$$

One should constrain the symmetry energy in multi-dimension parameter space

$$f = \frac{m}{k} [t_0 (2x_0 + 1) + t_3 (2x_3 + 1)] \frac{\rho}{2}$$

Model: transport code, ImQMD

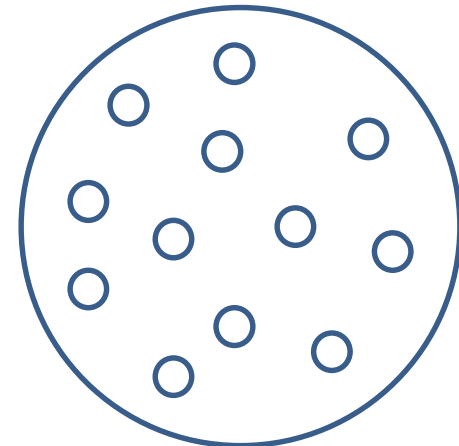
Ab-initio



$$\Psi(r_1, r_2, r_3, \dots, r_N)$$

$$\langle \Psi | H | \Psi \rangle = E$$

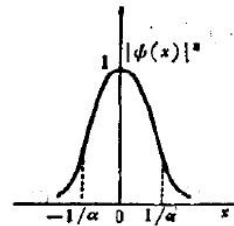
QMD approaches



Hartree approximation
Mean field level

$$\Psi \approx \Phi = \prod_i \phi_i$$

$$H \rightarrow H_{eff}$$



$$\frac{\partial \langle p_i \rangle}{\partial t} = - \langle \nabla_i V(R_1, \dots, R_N) \rangle \approx - \frac{\partial U(\bar{R}_1, \dots, \bar{R}_N)}{\partial \bar{R}_i}$$

$$\frac{d \langle R_i \rangle}{dt} = \frac{\langle p_i \rangle}{m}$$

We can use EDF in QMD approach

$R_{i0} = \langle R_i \rangle$ and $p_{i0} = \langle p_i \rangle$ are the centroid of wave packets in coordinate and momenta

Nucleon-nucleon collision integral part is simulated with cascade

- ImQMD with standard Skyrme interaction**

the potential energy U that includes the full Skyrme potential energy without the spin-orbit term:

$$U = U_\rho + U_{md} + U_{coul} \quad (2)$$

and U_{coul} is the Coulomb energy. The nuclear contributions are represented in a local form with

$$U_{\rho,md} = \int u_{\rho,md} d^3r$$

$$u_\rho = \frac{\alpha}{2} \frac{\rho^2}{\rho_0} + \frac{\beta}{\eta + 1} \frac{\rho^{\eta+1}}{\rho_0^\eta} + \frac{g_{sur}}{2\rho_0} (\nabla \rho)^2$$

$$+ \frac{g_{sur,iso}}{\rho_0} [\nabla(\rho_n - \rho_p)]^2$$

$$+ A_{sym} \rho^2 \delta^2 + B_{sym} \rho^{\eta+1} \delta^2$$

and

Skyrme type Momentum Dependent Interaction

$$\delta(r_1 - r_2)(p_1 - p_2)^2$$

Y.X. Zhang, M.B.Tsang, Z.X. Li, H.Liu, PLB732,186 (2014)

$$u_{md} = u_{md}(\rho\tau) + u_{md}(\rho_n\tau_n) + u_{md}(\rho_p\tau_p) \quad (5)$$

$$= C_0 \int d\vec{p} d\vec{p}' f(\vec{r}, \vec{p}) f(\vec{r}, \vec{p}') (\vec{p} - \vec{p}')^2$$

$$+ D_0 \left[\int d\vec{p} d\vec{p}' f_n(\vec{r}, \vec{p}) f_n(\vec{r}, \vec{p}') (\vec{p} - \vec{p}')^2 + \int d\vec{p} d\vec{p}' f_p(\vec{r}, \vec{p}) f_p(\vec{r}, \vec{p}') (\vec{p} - \vec{p}')^2 \right]$$

Interaction parameter \rightarrow Nuclear matter parameters in SHF

$$g_{\rho\tau} = \frac{3}{5} \left(\frac{m_0}{m_s^*} - 1 \right) \epsilon_F^0,$$

$$\gamma = \frac{K_0 + \frac{6}{5} \epsilon_F^0 - 10g_{\rho\tau}}{\frac{9}{5} \epsilon_F^0 - 6g_{\rho\tau} - 9E_0}$$

$$\beta = \frac{(\frac{1}{5} \epsilon_F^0 - \frac{2}{3} g_{\rho\tau} - E_0)(\gamma + 1)}{\gamma - 1}$$

$$\alpha = E_0 - \epsilon_F^0 - \frac{8}{3} g_{\rho\tau} - \beta$$

$$C_{sym} = -\frac{1}{24} \left(\frac{3\pi^2}{2} \right)^{2/3} (3\Theta_v - 2\Theta_s) \rho^{5/3}$$

$$B_{sym} = \frac{3S_0 - L - \frac{1}{3} \epsilon_F^0 + 2C_{sym}(m_s^*, m_v^*)}{-3\sigma}$$

$$A_{sym} = S_0 - \frac{1}{3} \epsilon_F^0 - B_{sym} - C_{sym}(m_s^*, m_v^*)$$

$$\Theta_s = \left(\frac{m}{m_s^*} - 1 \right) \frac{8\hbar^2}{m} \frac{1}{\rho_0}, \quad \Theta_v = \left(\frac{m}{m_v^*} - 1 \right) \frac{4\hbar^2}{m} \frac{1}{\rho_0}.$$

7 nuclear matter parameters!

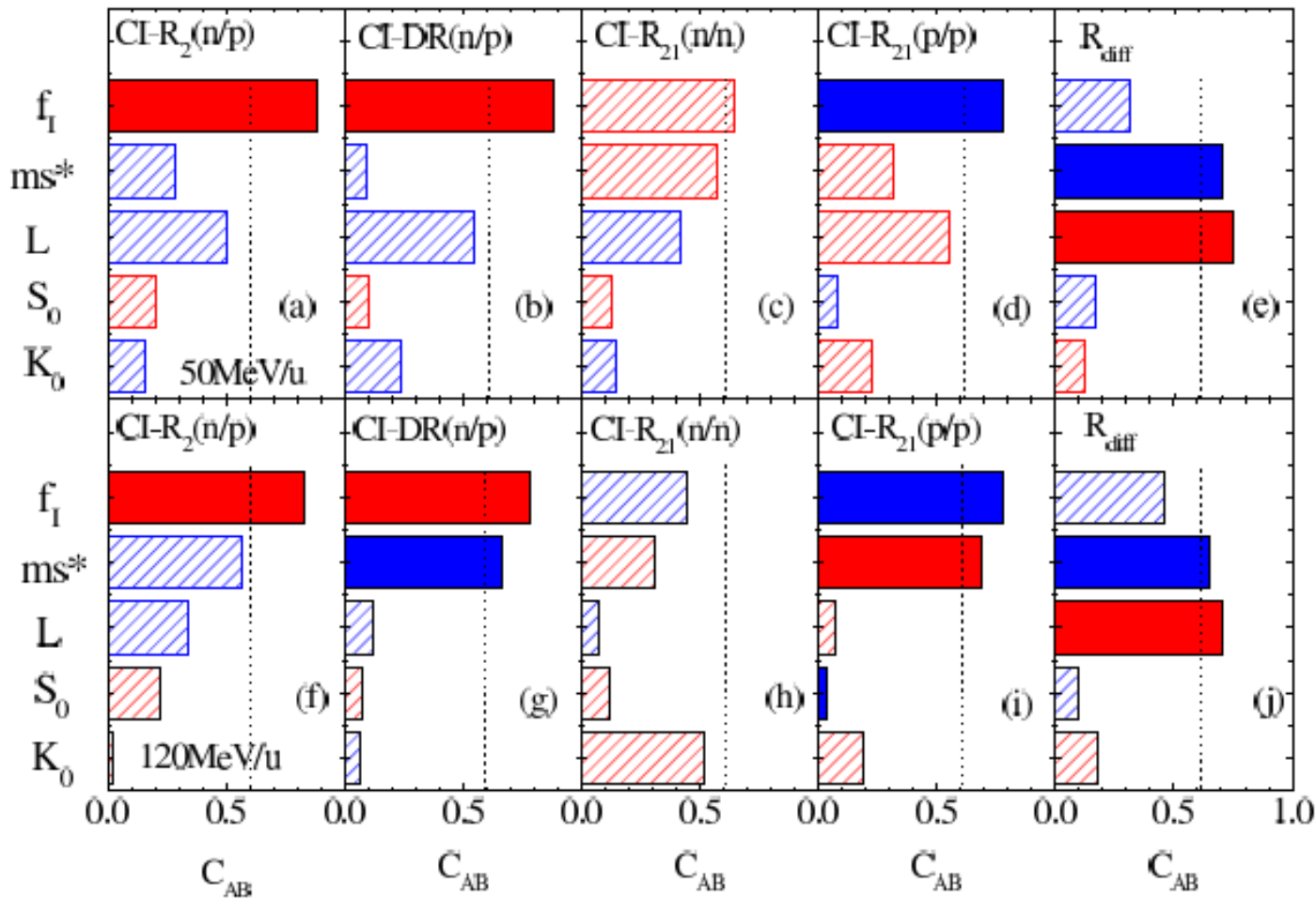
$$C_0 = \frac{1}{16\hbar^2} [t_1(2 + x_1) + t_2(2 + x_2)] = \frac{1}{16\hbar^2} \Theta_v$$

$$D_0 = \frac{1}{16\hbar^2} [t_2(2x_2 + 1) - t_1(2x_1 + 1)] = \frac{1}{16\hbar^2} (\Theta_s - 2\Theta_v)$$

Sensitivity analysis of model parameter
 $x=\{\}$ to $O=\{\}$

Sensitivity analysis C_{AB} of $x=\{\}$ and $O=\{\}$

Y.X.Zhang, M.B.Tsang, Z.X.Li, PLB749,262(2015)



Blue:
negative correlation

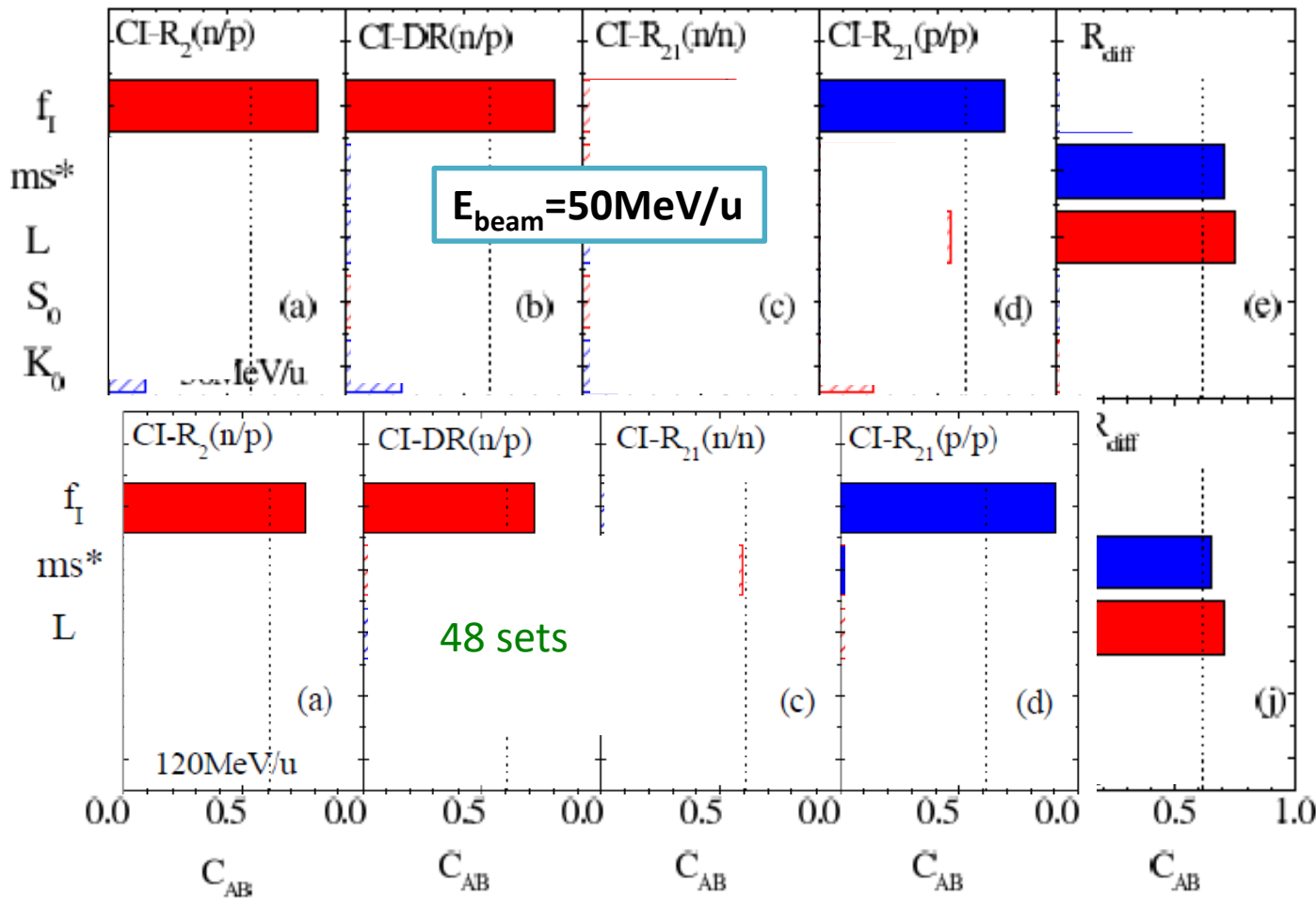
Red:
Positive correlation

The ratios are
constructed with
Ek>40MeV

- Ms^* also play important roles for isospin diffusion, and neutron to proton yield ratio observables at 120MeV/u.

Sensitivity analysis C_{AB} of $x=\{\}$ and $O=\{\}$

Y.X.Zhang, M.B.Tsang, Z.X.Li, PLB749,262 (2015)



Blue:
negative correlation

Red:
Positive correlation

The ratios are constructed with $E_k > 40 \text{ MeV}$

- Ms^* also play important roles for isospin diffusion, and neutron to proton yield ratio observables at 120 MeV/u , one can reasonable determine $\{K_0, S_0, L, ms^*, f_i\}$ by combination analysis.**

Bayesian inference on the $\{K_0, S_0, L, ms^*, f_I\}$ in multi-dimension parameter space for Skyrme parameter set

$$P(X|D) = \frac{P(D|X)P(X)}{P(D)}$$

- Prior information of X

$$P(X)$$

- Likelihood function

$$P(D|X) \quad L(x) \sim \exp\left(-\sum_a \frac{(y_{M,a}(x) - y_a^{exp})^2}{2\sigma_a^2}\right)$$

- Likelihood function

$$L(x) \sim \exp\left(-\sum_a \frac{(y_{M,a}(x) - y_a^{exp})^2}{2\sigma_a^2}\right)$$

7 Nuclear matter parameters:

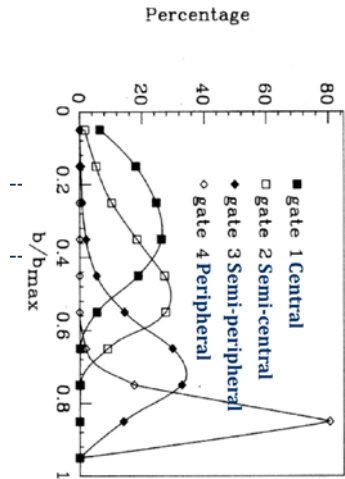
$\{\rho_0, E_0, K_0, S_0, L, m_s^*, m_v^*\}$

Use $\rho_0=0.16$, $E_0=-16\text{MeV}$

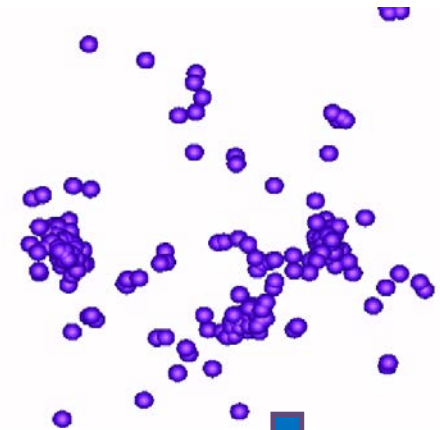
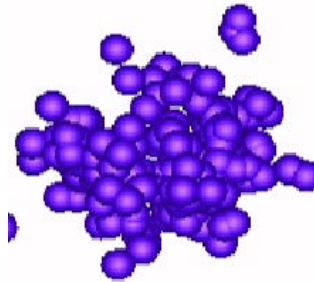
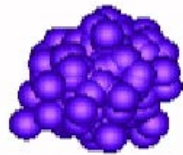
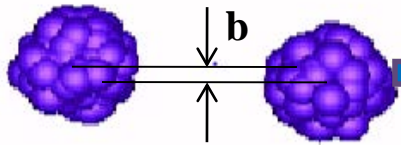
- $K_0=200-280\text{ MeV}$
- $S_0=25-35\text{ MeV}$
- $L=30-120\text{ MeV}$,
- $m_s^*/m=0.6-1.0$
- $f_i=-0.5-0.4$

120 set points (training points)

Considering the Impact parameter smearing



Estimate
impact parameter



M. B. Tsang et al. PLB 220, 492 (1989);

L. Phair et al. NPA 548, 489-509 (1992);

J.F. Lecomte et al. PLB 325, 317-321 (1994);

J. Lukasik et al. PRC 55, 1906 (1997);

T. X. Liu et al. PRC 86, 024605 (2012).

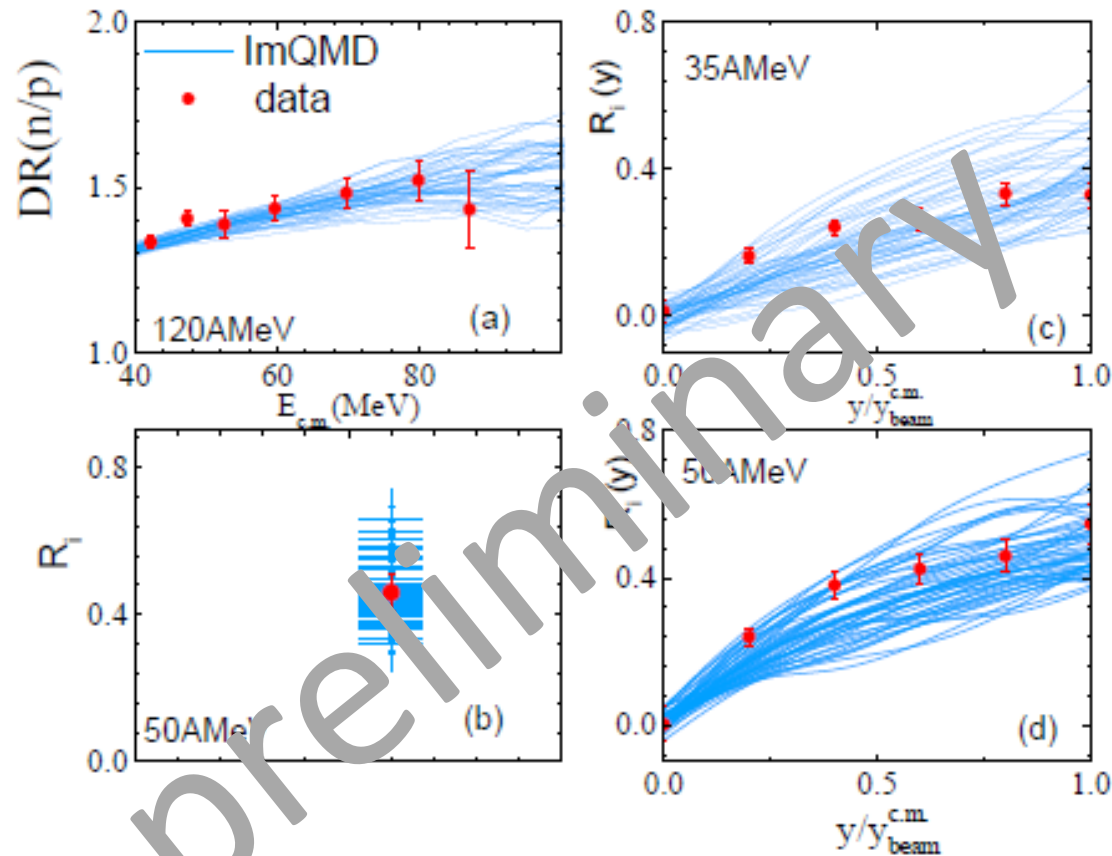
...

1. M , 2. E_{t12} , 3. θ_{flow} ...

Go back to the process

Observables

$$\{\rho_0, E_0, K_0, S_0, L, m_s^*, m_v^*\}$$



obtained with ImQMD, for 4 Obs from 88 sets

Data:

D.D.Coupland, et.al., PRC94,011601(2016)

Z.Y.Sun, et.al., PRC82, 051603(2010)

T.X.Liu, et.al., PRC76,034603(2007)

Tsang, et.al., PRL92,062701(2006)

4, Summary and outlook

1), A new version ImQMD can use the 'real' Skyrme EDF was developed

2), parameter correlation should be carefully considered in the reliable constraints on symmetry energy in HICs

3), Bayesian inference on Skyrme EDF from HIC together with neutron skin and NS properties is needed.

Collaborator:

M.B.Tsang (曾敏儿), NSCL/Michigan State University

Pierre Morfauc, NSCL/MSU

Hang Liu (刘航), Texas Advanced Computer center,
University of Texas,

Zhuxia Li (李祝霞), China Institute of Atomic Energy,

NSCC-TJ for calculations

Thanks for your attention!